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LA-10 Series 3

THE EXPONENTIAL SHOCK AND MARKACTION WAVE

The following curves give the results of the numerical solution of the equations of motion of the one dimensional shock and rarefaction waves derived by R. Serber. These waves occur at the interface of two materials, one of which (the core) is heated by an energy release, the spatial distribution of which is constant per gram, and which increases exponentially with time. Each of the two materials is assumed to have zero initial temperature and to obey the equation of state of a perfect gas. The calculations for the shock wave have been carried thru for two values of the ratio of specific heats, $\gamma = 1.4$ and 1.67. The refaction wave has been treated only for $\gamma = 1.4$.

sions increases exponentially, as do the scales of pressure and temperature. If the scales are so changed the graphs giving the spatial distribution of pressure, temperature, and displacement are invariant. The density as a function of the exponentially increasing scale of length is invariant. The scale of length which has been used is the co-moving coordinate, i.e. the initial distance from the interface of the material at the point considered, normalised to unity at the head of the wave. This length, denoted by x, has the range O to 1 for the shock wave and -1 to O for the rarefaction wave. For the rarefaction wave the letter \$\eta\$ has been used in place of \$\pi\$.

The actual pressure or temperature at any stage of the expansion can be determined by making use of the fact that the head of the rarefaction wave moves with sound velocity, [P/p] **. For a specified rate of increase of energy density, pressure, and temperature, and of position and velocity, etc., this condition determines the pressure and temperature at the head of the rarefaction wave in terms of its position and density. The pressure is continuous across the

SEE A REPORT OF R. SERBER TO AFPEAR LATER IN THIS SERIES, FOR FURTHER CLARIF CATION.







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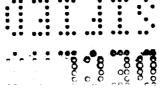
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interface. This determines the pressure, hence also the temperature, in the shock seve.

The character of the shock wave is determined by Ytamper alone. The character of the rarefaction wave is then determined by the interface boundary condition which equates the pressure and displacement on the two sides of the interface, hence also the invariant ratio, pressure over displacement squared. This condition depends on the two values of Y and on the ratio of initial densities, $\alpha = \rho_{tamper}/\rho_{core}$:

$$\frac{2\rho_t \quad P(0)}{(\gamma_t+1)P(1) \quad \varphi^*(+0)} = \frac{\rho_c \quad F(0)}{\gamma_c \quad P(-1) \quad \varphi^*(-0)}$$

In the curves given here y has been taken to be 1.4 in both materials.

The notation used in the graphs is as follows:

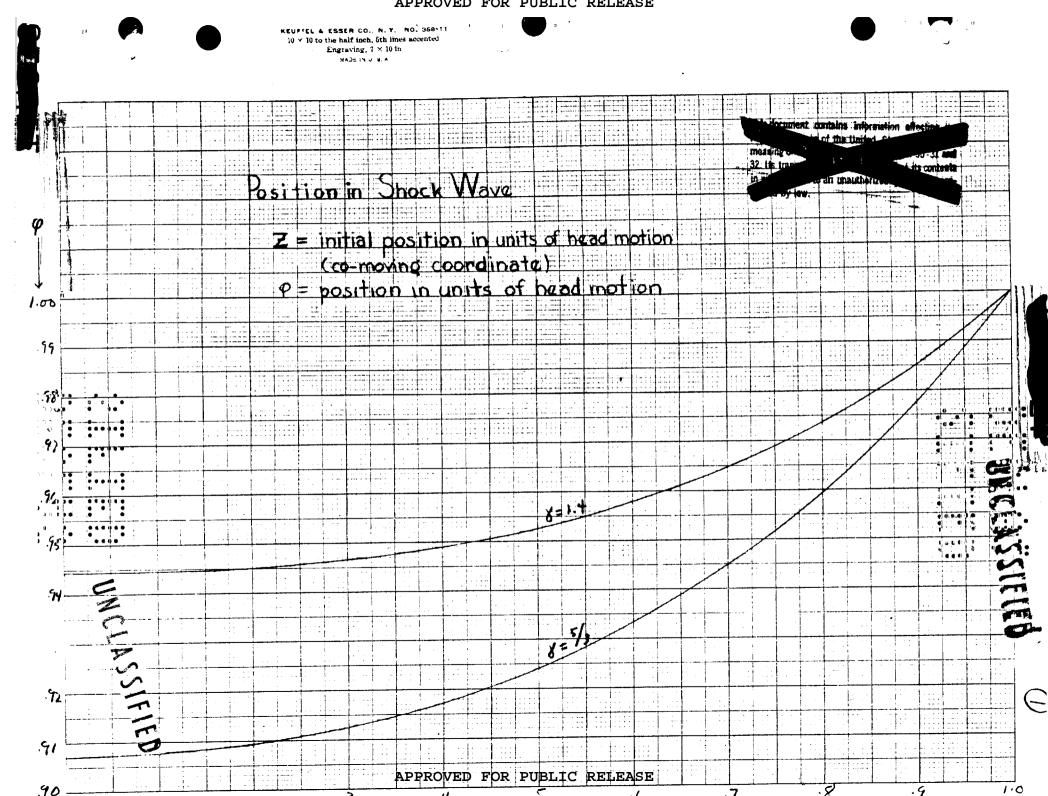
- ρ. initial density of core or tamper
- Pet/Pec
- n the initial position (wessured from the initial rosition of the interface) of the material of the rosition of the head of the water or ratefaction wave.
- = sctust rosition in units of the head position
- λ = the displacement, 🛻 🖔
- P = pressure in units of pressure at head of wave
- ρ = density in units of ρ_0
- $T = P/\rho$, temperature

It should be observed that the units in which z, , F, p, and T are measured are not neccessarily the same in the tamper as in the core. The relationship between the two sets of units is determined by the equality of pressure and position at the interface.

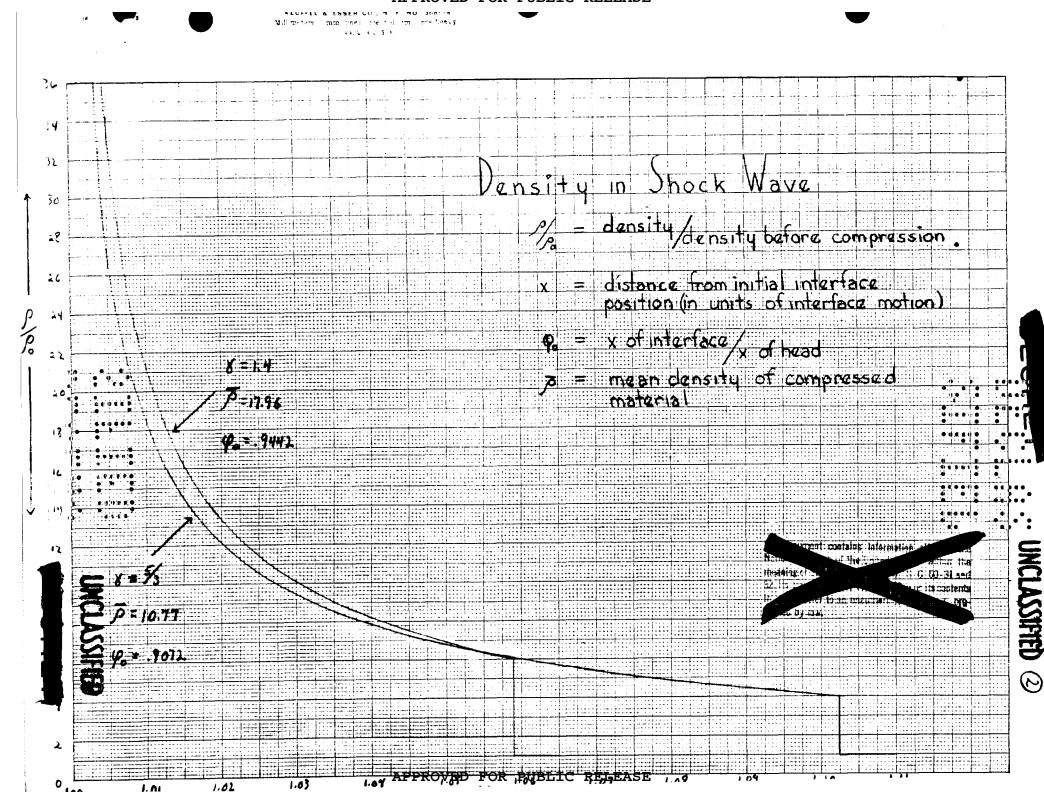
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